# Challenges of Solar Cycle Prediction Introduced by 'Rogue' Active Region Emergences <u>and</u> Meridional Inflow

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# Challenges of prediction

#### 1. Rogue BMRs

One single 'rogue' BMR, as it can have a flux content commensurate to the maximum of the polar cap flux, has a huge impact on the subsequent cycle. Due to this fact, it is very important to treat these extreme events in prediction.

- 2. Meridional inflow Potential non-linear mechanism
- **3. Torsional oscillation** next step

### The kinematic "2×2D" Babcock-Feighton dynamo model

Lemerle et al., 2015, ApJ 801; Lemerle & Charbonneau, 2017, ApJ 834





# Solar-like Solution - prediction





# Meridional inflow $\rightarrow$ flow perturbation



- The inflows are caused by flows towards individual active regions due to the local cooling effect of the magnetic field.
- Surface effect, decreases quickly with depth.
- According to measurements, the inflow strength is about 3-5 m/s on average
- H. C. Spruit 2003 SoPh 213

Model of inflows towards an activity complex formed by several emerged BMRs. The colour scale encodes the strength of the magnetic field B<sub>r</sub> D. Martin-Belda and R. H. Cameron 2017 A&A **597** 

# Meridional inflow $\rightarrow$ flow perturbation



Helioseismic measurements of the meridional flow during cycle 23: the perturbation caused by the inflows is the strongest at cycle maximum. I. Gonzalez Hernandez et al. 2010 *ApJ* **713** 

The inflows represent a **potential non-linear mechanism** capable of **limiting the build up of the polar fields** and **contributing to the modulation of the solar cycle**.

R. H. Cameron and M. Schüssler 2010 ApJ 720



# Models of the meridional inflow

#### J. Jiang et al. 2010 ApJ 717

Perturbation in the meridional flow in their Surface Flux Transport model. The inflow belt moves together with the activity belt, fix width and flowspeed. Result: **reduction of polar fields by ~18%** 

$$v(\lambda, t) = \begin{cases} v_{\rm m} \sin(2.4\lambda) + \Delta v(\lambda, t) & \text{for } |\lambda| \leq 75^{\circ} \\ 0 & \text{otherwise,} \end{cases}$$

where  $v_{\rm m} = 11 \text{ m s}^{-1}$  and

$$\Delta v(\lambda, t) = \begin{cases} v_0 \sin \left[ (\lambda - \lambda_c(t)) / \Delta \lambda_v \right] & \text{for } -180^\circ \\ \leqslant (\lambda - \lambda_c(t)) / \Delta \lambda_v < 180^\circ \\ 0 & \text{otherwise.} \end{cases}$$



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R. H. Cameron and M. Schüssler 2010 *ApJ* **720** Magnetic field dependent flow speed. ← **The flow should depend on the strength of the cycle.** 

$$v(\lambda, t) = c_0 \left\{ \frac{d\langle |B| \rangle_{\phi}(\lambda, t)}{d\lambda} \right\}$$



## Meridional inflow in the 2×2D model

$$\Delta v(\lambda, t) = \begin{cases} v_0 \sin \left[ (\lambda - \lambda_c(t)) / \Delta \lambda_v \right] & \text{for } -180^\circ & \leq (\lambda - \lambda_c(t)) / \Delta \lambda_v < 180^\circ \\ 0 & \text{otherwise.} \end{cases}$$

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$$\Delta v(t, F) = v_0(t, F) \sin\left(\frac{\lambda - \lambda_c(t, F)}{\Delta \lambda(t, F)}\right), \quad \text{for: } -180^\circ \le \frac{\lambda - \lambda_c(t, F)}{\Delta \lambda(t, F)} < 180^\circ$$

Based on the concept of Jiang et al., BUT:

- Flow parameters defined with the parameters of the emerged active regions.
- Calculated separately for the two hemispheres asymmetry
- Updated at each internal dynamo timestep

## Meridional inflow in the 2×2D model

$$\Delta v(t, F) = v_0(t, F) \cdot \sin\left(\frac{\lambda - \lambda_c(t, F)}{\Delta \lambda(t, F)}\right), \quad \text{for: } -180^\circ \le \frac{\lambda - \lambda_c(t, F)}{\Delta \lambda(t, F)} < 180^\circ$$

$$v_0(t, F) = v_{00} \cdot \arctan\left(\frac{\Sigma F}{F_{00}}\right)$$

#### **Flow velocity**

 $v_{00} = 5 \text{ m/s} - \text{observations} (5.25 \text{ m/s}; 5.50 \text{ m/s}; 7.50 \text{ m/s})$   $F_{00} = 4.99 \, 10^{21} \text{ Mx} \text{ (average unsigned flux)}$  $F_i - \text{flux of BMRs emerged during an internal dynamo timestep}$ 

$$\lambda_c(\mathbf{t}, \mathbf{F}) = \frac{\Sigma(\lambda_i \cdot F_i)}{\Sigma F_i}$$

Flux-weighted centroid of the inflow belt  $\lambda_i$  –midpoints of BMRs

$$\Delta \lambda(\mathbf{t}, \mathbf{F}) = \frac{1}{2} \sqrt{\frac{1}{n} \left(\lambda_i - \lambda_c(\mathbf{t}, \mathbf{F})\right)^2}$$

Width of the inflow belt n – number of emerged BMRs on the North / South



#### **Reduced stochasticity**

#### **Full stochasticity**

(tilt angle varies according to Joy's law; nonstochastic separation)



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# Conclusions

Flux dependent meridional inflow is added to the 2×2D Babcock-Leighton solar dynamo model.

#### Results:

- The inflow reduces the dipole moment by ~10% → reduction in the cycle amplitudes, too (cycle prediction!)
- Cross-equatorial flows → stronger hemispheric coupling (asymmetry!)

#### Next step:

- Link cross-equatorial flows to hemispheric asymmetry
- More case studies to investigate the impact of meridional inflows on cycles with different amplitude (different effects are proposed)

# Thank you for your attention!

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# Models of the meridional inflow

- R. H. Cameron and M. Schüssler 2012 A&A 548
- Cross-equatorial flows:
- BMRs close to the equator: enhanced flux transport between hemispheres -> it strengthens the dipole moment (in case of <u>weak cycles</u>)
- In general:

converging flows decrease the tilt angle of the BMRs -> lower contribution to the dipole moment (in case of <u>strong cycles</u>, where the activity belts are further from the equator, this effect will **weaken the dipole moment**)

#### Nonlinear feedback!



#### **Reduced stochasticity**

#### **Full stochasticity**

(tilt angle varies according to Joy's law; non-stochastic

separation)





Full simulation with flux dependent inflow and stochastic tilt & separation.

- Cycle-to-cycle variation
- Hemispheric asymmetry
- Cross-hemispheric flows

## histogram of $f_{f/i}(\theta, t) \cdot \delta D_{BMR}$



# Method



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